



CoF3 FeF3 40 K CoF3 Co DM Co

Exploring novel magnetoelastic coupling via polyhedral or molecular rotation in unusual quantum spin systems, we studied perovskite-type TF3(T=Sc, Ti, Cr, Mn, Fe, Co), superoxide KO2, and etc. The crystal and magnetic structural analysis is going well. E. O. Wollan et al.(Phys. Rev. 112 1132) reported magnetic structure of CoF3 and FeF3. However, (101) and (003) hexagonal peaks were not distinguished well due to poor resolution of neutron powder diffractometer. By using high resolution neutron powder diffractometer, we confirm the magnetic structure at room temperature. The Fe antiferromagnetic moment is inside hexagonal plane while Co antiferromagnetic moment is parallel with c-axis. DM interaction induces spin canting in the case of FeF3 that causes weak-ferromagnetism. In addition, we find low-temperature weak-ferromagnetic transition in CoF3. This result suggest the Co spin canting will induces weak-ferromagnetism by DN interaction below 40 K.

Magnetism

Perovskite Fluoride Superoxide Neutron diffraction Magnetism Crystallography Phase tran sition

When magnetic moment ordered in condensed matter, it interplays with another degree of freedom such as the lattice. It is called magnetoelastic coupling which type is strictly followed by crystal and magnetic symmetry.

MnO, GO, and N<sub>O</sub> are standard antiferromagnetic materials which is extensively studied for last several decades. However, detailed crystal and magnetic structure, phase transition mechanism are so debates. Recently, we found spin-direction-dependent magnetoelastic coupling in transition metal monoxide [1]. The exact spin direction or spin symmetry is hidden mechanism to distinguish different magnetic phase transition. The MnO is gamma1 magnetic structure with discontinuous phase transition. Their spin direction is fixed symmetrically. On the other hand, NiO and CoO are gamma2 magnetic structure with continuous phase transition. The spin direction can be rotated. In addition, Gamma2 magnetic structure show unusual linear-cubic magnetoelastic coupling forbidden by time-reversal symmetry.

These results inspire me to explore unusual magnetoelastic coupling forbidden by time-reversal symmetry. It is related with the issue of general magnetic symmetry beyond typical time-reversal symmetry.

The double antisymmetry space group introduces rotation-time reversal symmetry [2]. If both rotational axial vector and magnetic axial vector can be reversed simultaneously by rotation-time reversal symmetry, it can allow unusual linear-cubic magnetoelastic coupling forbidden by time-reversal symmetry alone. Because of the edge sharing octahedral of transition metal monoxide, double antisymmetry space group is difficult to apply the case of MnO, GoO, and NO directly. Thus, we interest the magnetic system which have polyhedral rotation and magnetic moments. The candidates are transition metal trifluoride, KQ<sub>2</sub>, and etc.

The samples are prepared from company or

collaborators. We characterize crystal structure by XRD and magnetic properties by SQUID. We employ time-of-flight neutron powder diffractometer, which is SuperHRPD beamline in MLF, J-PARC, to study detailed crystal and magnetic structure simultaneously. The Rietveld analysis is carried out by Z-Rietveld and Fullprof. For advance crystallography and group theory, we used SARAh, bilbaocrystallographic server and ISOTROPY.

The several experimental results are still under analysis. Those results will be published as soon as possible after summary. Here, I introduce briefly the recent publication of  $C\Phi_3$  and FeF<sub>3</sub> study [3].



Fig.1 Gystal structure  $(R3c, aaa)$  of  $C\circ F_3$  and FeF<sub>3</sub>.



Fig. 2 Time-of-flight neutron diffraction of CoF3 and FeF3 at room temperature.



Fig. 3 Magnetic susceptibility measurement of  $C\overline{or}_3$  and  $F\overline{er}_3$ .



 $H_{DM} = \sum_{i < i} D_{ii} (S_i \times S_i)$ 

Fig. 4 Weak-ferromagnetism can be induced by DM interaction. At room temperature, weak-ferromagnetism is allowed in FeF<sub>3</sub> magnetic structure.

Cobalt trifluoride  $C\circ F_3$  and Iron trifluoride Fe $F_3$  have **R**3c space group which is described by Glazer tilting system a-a-a-. From simple cubic perovskite *PinBríp*a<sup>0</sup>a<sup>0</sup>, The octahedral rotates along [100], [010], and [001] pseudocubic direction with antiphase rotation of nearby octahedral. Figure 1  $d$ isplay CoF<sub>3</sub> and FeF<sub>3</sub> crystal structure in hexagonal setting.

These CoF3 and FeF3 crystal and magnetic structure were investigated by early angle-dispersive neutron powder diffractometer[4]. However, it couldn't distinguish 101, 003 magnetic peaks clearly. Thus, proposed magnetic structure are not confirmed completely. As shown in Figure 2, we distinguish 101, 003 magnetic peaks unambiguously that confirm previously proposed magnetic models. At room temperature, Co spins are parallel with hexagonal c-axis while Fe spins are perpendicular to c-axis.

In Fig. 3, magnetic susceptibility measurement shows additional weak-ferrimagnetic transition of  $C\Phi_3$  at 40 K which is below well-known antiferromagnetic transition 460 K.

Dzyaloshinsky-Moriya(DM) interaction is antisymmetric superexchange interaction which microscopic origin is spin-orbit coupling. DM interaction induces spin canting and cause weak-ferromagnetism. The DM interaction is acting through magnetic atom-ligand-magnetic atom path when there is no inversion symmetry at ligand atomic position. The octahedral rotations shift ligand atoms from inversion symmetry of ideal cubic perovskite. Thus, octahedral rotations are necessary condition for DM interaction.

Since the uncancelled DM vector is parallel with c-axis, DM interaction can induce weak-ferromagnetism such as  $FeF<sub>3</sub>$ case. Fe spins are perpendicular to c-axis. Canted Fe spin inside hexagonal plane can reduce total energy through DM interactions while Co spins are parallel with c-axis which has no lower energy via DM interaction. Therefore, only FeF<sub>3</sub> has weak-ferromagnetism at room temperature.

However, as shown in Fig. 3, we find additional weak-ferromagnetism in  $C\delta F_3$ unexpectedly. The room-temperature  $C\sigma F_3$ magnetic structure can't cause weak-ferromagnetism. In-plan Co spin component is necessary for weak-ferromagnetism. From this antiferromagnetic spin direction and DM interaction, we expect that Co spin will be rotated. After that, in-plan Co spin component will induce weak-ferromagnetism. We carry out further structural study to clarify this weak-ferromagnetism mechanism. It is under analysis.

S. Lee et al., Phys. Rev. B 83, 065539 (2016)

V. Gopalan and D. B. Litvin, Nat. Mater. 10, 376 (2011)

S. Lee et al, Physica B: Condensed Matter (2017)

DOI: 10.1016/j.physb.2017.11.082

E. O. Wollan et  $d.$ , Phys. Rev. 1122 1132 (1958)

Sanghyun Lee, Shuki Torii, Yoshihisa Ishikawa, Masao Yonemura, Taketo Moyoshi, Takashi Kamiyama, Physica B: Condensed Matter (2017)

DOI: 10.1016/j.physb.2017.11.082

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16 th Korea-Japan meeting on Neutron Science, Kashiwa in Japan (Jan. 8-10, 2018)

24 th Congress and General Assembly of the international Union of Crystallography, Hyderabad in India (21-28, August 2017)

MRS-Ina C&C 2017, Yogyakarta, Indonesia (8-12, October, 2017)

Theme Meeting on Neutron Scattering,

Mumbai in India (19, Aug, 2017)

International Conference on Neutron Scattering, Deajeon in Korea, (July 9-13, 2017)

2016

Tsukuba in Japan (March 14-15, 2017) 14<sup>th</sup> Conference of the Asian Crystallographic Association, Hanoi in

Vietnam (4-7 December 2016)

EMN Dalian Meeting, Dalian in China (25-29, July, 2016)

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